

Reduced order model approach for imaging with waves

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Velocity estimation problem

- *Direct problem:* Given the velocity map $c = (c(x))_{x \in \Omega}$ compute the wavefield solution of the wave equation

$$[\partial_t^2 - c^2(x)\Delta]p^{(s)}(t, x) = f(t)\delta(x - x_s), \quad t \in \mathbb{R}, x \in \Omega,$$

starting from $p^{(s)}(t, x) = 0, t \ll 0$, + boundary conditions at $\partial\Omega$.
At the locations of the receivers:

$$d_{r,s}(t) = p^{(s)}(t, x_r), \quad r, s = 1, \dots, N$$

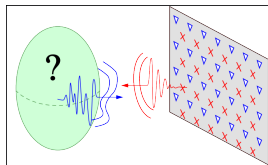
\hookrightarrow forward map

$$\mathcal{D} : c \mapsto \mathbf{d}$$

where $\mathbf{d} = ((d_{r,s}(t)))_{r,s=1}^N_{t \in [t_{\min}, t_{\max}]}$, is the array response matrix.

- *Inverse problem:*

Given the time-resolved measurements \mathbf{d} ,
determine the velocity map c .



Full Waveform Inversion (FWI)

- FWI fits data with its model prediction in L^2 -norm:

$$\hat{c} = \operatorname{argmin}_c \{ \mathcal{O}_{FWI}[c] + \operatorname{Reg}[c] \},$$

$$\mathcal{O}_{FWI}[c] = \|\mathbf{d}_{meas} - \mathcal{D}[c]\|_2^2 = \sum_{r,s=1}^N \int_{t_{\min}}^{t_{\max}} |d_{meas}(t)_{r,s} - \mathcal{D}[c](t)_{r,s}|^2 dt$$

Cf [Virieux and Operto 2009].

- Bayesian interpretation: Maximum A Posteriori (with Gaussian additive noise).
- The objective function $\mathcal{O}_{FWI}[c]$ is not convex in c .
 \hookrightarrow optimization needs hard to get good initial guess.

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- FWI fits data with its model prediction in L^2 -norm:

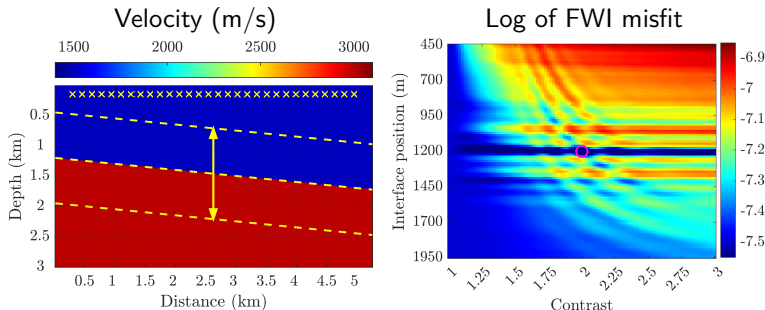
$$\hat{c} = \underset{c}{\operatorname{argmin}} \{ \mathcal{O}_{FWI}[c] + \operatorname{Reg}[c] \},$$

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- Bayesian interpretation: Maximum A Posteriori (with Gaussian additive noise).
- The objective function $\mathcal{O}_{FWI}[c]$ is not convex in c .
 \hookrightarrow optimization needs hard to get good initial guess.
- Layer stripping: Proceed hierarchically from the shallow part to the deep part [Wang et al. 2009]
- Frequency hopping: Successive inversion of subdata sets of increasing high-frequency content [Bunks et al. 1995]
- Optimal transport: Wasserstein distance instead of L^2 -norm [Engquist et al. 2016, Yang et al. 2018, Métivier et al. 2019]

Topography of the FWI objective function



- Probing pulse is a modulated Gaussian pulse with central frequency 6Hz and bandwidth 4Hz ($\lambda \simeq 300\text{m}$ at 10Hz).
- $N = 30$ sensors and $N_t = 39$ time samples at interval $\tau = 0.0435\text{s}$.
- Search velocity has two parameters: the bottom velocity and depth of the interface (the angle and top velocity are known).
- Objective function:

$$\mathcal{O}_{FWI}[c] = \|\mathbf{d}_{meas} - \mathcal{D}[c]\|_2^2$$

↪ Cycle skipping issues (many local minima).

Main objective

- Our main objective: find a convex formulation of FWI.

Proposed approach: Find a nonlinear mapping \mathcal{R} : data $\mathbf{d} \mapsto$ reduced order model (ROM) of wave operator \mathbf{A}^{rom} (matrix) such that:

- ROM is computed with efficient numerical linear algebra tools in non-iterative fashion.
- Minimization of ROM misfit is better for velocity estimation.

We can think of the data to ROM mapping \mathcal{R} as a preconditioner of the forward mapping \mathcal{D}

$$\mathbf{c} \xrightarrow{\mathcal{D}} \mathbf{d} \xrightarrow{\mathcal{R}} \mathbf{A}^{rom}$$

because the composition $\mathcal{R} \circ \mathcal{D}$, which gives $\mathbf{A}^{rom} = \mathcal{R}(\mathcal{D}[\mathbf{c}])$, is easier to “invert”.

Towards the ROM objective function

- *Ideal objective function 1:*

$$\mathcal{O}[c] = \|c - c^{meas}\|^2$$

but c^{meas} is not observed (i.e., cannot be extracted from \mathbf{d}_{meas}) !

Towards the ROM objective function

- Let us consider the wave operator

$$\mathcal{A}[c] = -c(x)\Delta[c(x) \cdot]$$

- Galerkin method to approximate the operator \mathcal{A} by a matrix:
 - consider a space of (piecewise polynomial) functions with basis $(\Psi_l(x))_{l=1}^L$,
 - consider the row vector field $\Psi(x) = (\Psi_1(x), \dots, \Psi_L(x))$ and define:

$$\mathbf{A}^\Psi = \int_{\Omega} dx \Psi(x)^T \mathcal{A} \Psi(x) \in \mathbb{R}^{L \times L}$$

- *Ideal objective function 2:*

$$\mathcal{O}[c] = \|\mathbf{A}^\Psi[c] - \mathbf{A}^{\Psi, meas}\|_2^2$$

but $\mathbf{A}^{\Psi, meas}$ is not observed !

The ROM matrix

- Our Galerkin approximation space:
 - consider a time discretization $\{t_j = j\tau\}_{0 \leq j \leq N_t}$ with uniform stepping τ ,
 - gather the waves $p^{(s)}(t, x)$ evaluated at $t = t_j$ for all the N sources:

$$\mathbf{p}_j(x) = \left(p^{(1)}(t_j, x), \dots, p^{(N)}(t_j, x) \right), \quad x \in \Omega.$$

(note: apply first a linear preprocessing).

- organize the first N_t snapshots in the NN_t dimensional row vector field:

$$\mathbf{U}(x) = (\mathbf{p}_0(x), \dots, \mathbf{p}_{N_t-1}(x)), \quad x \in \Omega.$$

- apply Gram-Schmidt orthogonalization onto $\mathbf{U}(x) = \mathbf{V}(x)\mathbf{R}$.

- Define ROM matrix:

$$\mathbf{A}^{rom} = \int_{\Omega} dx \mathbf{V}(x)^T \mathcal{A} \mathbf{V}(x) \in \mathbb{R}^{NN_t \times NN_t}.$$

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$$U(x) = (p_0(x), \dots, p_{N_t-1}(x)), \quad x \in \Omega.$$

- apply Gram-Schmidt orthogonalization onto $U(x) = V(x)R$.

- Define ROM matrix:

$$\mathbf{A}^{rom} = \int_{\Omega} dx V(x)^T \mathcal{A} V(x) \in \mathbb{R}^{NN_t \times NN_t}.$$

- *Ideal objective function 3:*

$$\mathcal{O}[c] = \|\mathbf{A}^{rom}[c] - \mathbf{A}^{rom, meas}\|_2^2$$

but $\mathbf{A}^{rom, meas}$ is not observed (neither \mathcal{A} nor $V(x)$ is observed) !

The ROM matrix

- Our Galerkin approximation space:
 - consider a time discretization $\{t_j = j\tau\}_{0 \leq j \leq N_t}$ with uniform stepping τ ,
 - gather the waves $p^{(s)}(t, x)$ evaluated at $t = t_j$ for all the N sources:

$$\mathbf{p}_j(x) = \left(p^{(1)}(t_j, x), \dots, p^{(N)}(t_j, x) \right), \quad x \in \Omega.$$

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- organize the first N_t snapshots in the NN_t dimensional row vector field:

$$\mathbf{U}(x) = (\mathbf{p}_0(x), \dots, \mathbf{p}_{N_t-1}(x)), \quad x \in \Omega.$$

- apply Gram-Schmidt orthogonalization onto $\mathbf{U}(x) = \mathbf{V}(x)\mathbf{R}$.

- Define ROM matrix:

$$\mathbf{A}^{rom} = \int_{\Omega} dx \mathbf{V}(x)^T \mathcal{A} \mathbf{V}(x) \in \mathbb{R}^{NN_t \times NN_t}.$$

- **Main result:** *The ROM matrix \mathbf{A}^{rom} can be extracted from the measurements \mathbf{d} , without knowing \mathcal{A} nor $\mathbf{V}(x)$.*

$\hookrightarrow \mathcal{O}_{ROM}[c] = \|\mathbf{A}^{rom}[c] - \mathbf{A}^{rom, meas}\|_2^2$ is a legitimate objective function.

First step: *Linear preprocessing.*

- Define the new data matrix $\mathbf{D}(t)$:

$$\mathbf{D}(t) = \mathbf{d}^f(t) + \mathbf{d}^f(-t), \quad \text{with } \mathbf{d}^f(t) = -f'(-t) *_t \mathbf{d}(t).$$

Second step: *Expression of the new data entries as wave correlations.*

- Introduce the solution $u^{(s)}(t, x)$ of the homogeneous wave equation

$$(\partial_t^2 + \mathcal{A})u^{(s)}(t, x) = 0, \quad t > 0, \quad x \in \Omega,$$

with boundary conditions on $\partial\Omega$, with initial state

$$u^{(s)}(0, x) = u_0^{(s)}(x) = \left| \hat{f}(\sqrt{\mathcal{A}}) \right| \delta(x - x_s), \quad \partial_t u^{(s)}(0, x) = 0.$$

It has the form

$$u^{(s)}(t, x) = \cos(t\sqrt{\mathcal{A}})u_0^{(s)}(x).$$

→ The entries of $\mathbf{D}(t)$ can be expressed as wave correlations:

$$D_{r,s}(t) = \int_{\Omega} dx u_0^{(r)}(x) u^{(s)}(t, x).$$

Third step: *Definition of the ROM.*

Let $\tau > 0$ be fixed.

- Gather the snapshots for all the N sources in the row vector fields

$$\mathbf{u}_j(x) = \left(u^{(1)}(j\tau, x), \dots, u^{(N)}(j\tau, x) \right), \quad 0 \leq j \leq N_t.$$

- Organize the first N_t snapshots in the NN_t dimensional row vector field:

$$\mathbf{U}(x) = (\mathbf{u}_0(x), \dots, \mathbf{u}_{N_t-1}(x)), \quad x \in \Omega.$$

- Apply Gram-Schmidt orthogonalization onto $\mathbf{U}(x) = \mathbf{V}(x)\mathbf{R}$.
(note: we have $\int_{\Omega} dx \mathbf{V}(x)^T \mathbf{V}(x) = \mathbf{I}_{NN_t}$).

- Define

$$\mathbf{A}^{rom} = \int_{\Omega} dx \mathbf{V}(x)^T \mathcal{A} \mathbf{V}(x)$$

Fourth step: *Expression of the ROM in terms of mass and stiffness.*

- Define the $NN_t \times NN_t$ “mass” and “stiffness” matrices:

$$\mathbf{M} = \int_{\Omega} dx \mathbf{U}^T(x) \mathbf{U}(x), \quad \mathbf{S} = \int_{\Omega} dx \mathbf{U}^T(x) \mathcal{A} \mathbf{U}(x)$$

- Since $\mathbf{U}(x) = \mathbf{V}(x) \mathbf{R}$, we get

$$\begin{aligned} \mathbf{M} &= \mathbf{R}^T \int_{\Omega} dx \mathbf{V}^T(x) \mathbf{V}(x) \mathbf{R} \\ &= \mathbf{R}^T \mathbf{R} \end{aligned}$$

and

$$\begin{aligned} \mathbf{A}^{rom} &= \int_{\Omega} dx \mathbf{V}(x)^T \mathcal{A} \mathbf{V}(x) = \mathbf{R}^{-T} \int_{\Omega} dx \mathbf{U}(x)^T \mathcal{A} \mathbf{U}(x) \mathbf{R} \\ &= \mathbf{R}^{-T} \mathbf{S} \mathbf{R} \end{aligned}$$

↪ \mathbf{A}^{rom} can be expressed in terms of \mathbf{M} and \mathbf{S} .

Fifth step: *Expression of the ROM in terms of data.*

The $N \times N$ blocks of the mass matrix \mathbf{M} are

$$\begin{aligned}\mathbf{M}_{i,j} &= \langle \mathbf{u}_i, \mathbf{u}_j \rangle_{L^2(\Omega)} = \langle \cos(i\tau\sqrt{\mathcal{A}})\mathbf{u}_0, \cos(j\tau\sqrt{\mathcal{A}})\mathbf{u}_0 \rangle_{L^2(\Omega)} \\ &= \langle \mathbf{u}_0, \cos(i\tau\sqrt{\mathcal{A}})\cos(j\tau\sqrt{\mathcal{A}})\mathbf{u}_0 \rangle_{L^2(\Omega)} \\ &= \frac{1}{2} \langle \mathbf{u}_0, [\cos((i+j)\tau\sqrt{\mathcal{A}}) + \cos(|i-j|\tau\sqrt{\mathcal{A}})]\mathbf{u}_0 \rangle_{L^2(\Omega)} \\ &= \frac{1}{2} \langle \mathbf{u}_0, \mathbf{u}_{i+j} + \mathbf{u}_{|i-j|} \rangle_{L^2(\Omega)} \\ &= \frac{1}{2} (\mathbf{D}_{i+j} + \mathbf{D}_{|i-j|}), \quad 0 \leq i, j \leq N_t - 1.\end{aligned}$$

Idem for the stiffness matrix \mathbf{S} .

↪ \mathbf{M} and \mathbf{S} can be expressed in terms of the data \mathbf{D} .

Algorithm for data-driven ROM matrix

Input: The matrices $\mathbf{d}(t) = (d_{r,s}(t))_{r,s=1}^N$ of measurements.

1. Compute $d_{r,s}^f(t) = -f'(-t) *_t d_{r,s}(t)$ and

$$\mathbf{D}_j = \mathbf{d}^f(j\tau) + \mathbf{d}^f(-j\tau), \quad 0 \leq j \leq 2N_t - 2.$$

2. Compute $\ddot{\mathbf{D}}_j = \ddot{\mathbf{f}}(j\tau) + \ddot{\mathbf{f}}(-j\tau)$, for $j = 0, \dots, 2N_t - 2$ with $\ddot{f}_{r,s}(t) = f_{r,s}''(t)$ using, e.g., the Fourier transform.

3. Calculate $\mathbf{M}, \mathbf{S} \in \mathbb{R}^{NN_t \times NN_t}$ with the block entries

$$\mathbf{M}_{i,j} = \frac{1}{2}(\mathbf{D}_{i+j} + \mathbf{D}_{|i-j|}) \in \mathbb{R}^{N \times N},$$

$$\mathbf{S}_{i,j} = -\frac{1}{2}(\ddot{\mathbf{D}}_{i+j} + \ddot{\mathbf{D}}_{|i-j|}) \in \mathbb{R}^{N \times N},$$

for $0 \leq i, j \leq N_t - 1$.

4. Perform block Cholesky factorization $\mathbf{M} = \mathbf{R}^T \mathbf{R}$.

Output: $\mathbf{A}^{rom} = \mathbf{R}^{-T} \mathbf{S} \mathbf{R}^{-1}$.

ROM objective function

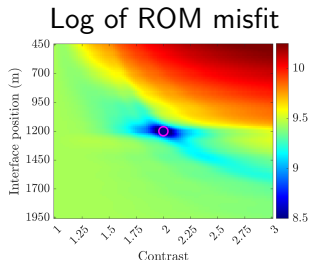
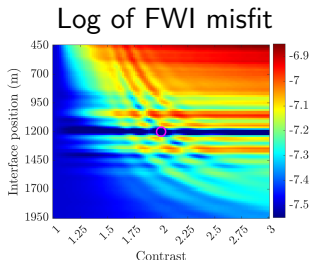
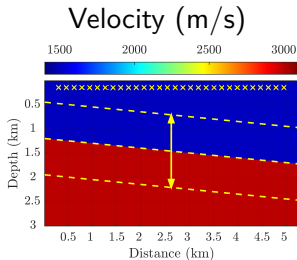
- ROM misfit function:

$$\mathcal{O}_{ROM}[c] = \|\mathbf{A}^{rom}[c] - \mathbf{A}^{rom, meas}\|_2^2$$

where $\mathbf{A}^{rom}[c]$ is computed from $\mathcal{D}[c]$ and $\mathbf{A}^{rom, meas}$ is computed from \mathbf{d}_{meas} .

- For a rich enough space of snapshots, the ROM matrix \mathbf{A}^{rom} contains roughly the same information as $\mathcal{A} = -c(x)\Delta[c(x) \cdot]$.
 \hookrightarrow The ROM misfit function should have nice convexity properties.
- Conjecture: “rich enough” means for sensors all around the domain of interest, separated by roughly half a wavelength, for time sampling satisfying the Nyquist criterium.
 \hookrightarrow Conjecture proved only in special situations.

Topographies of the FWI and ROM objective functions



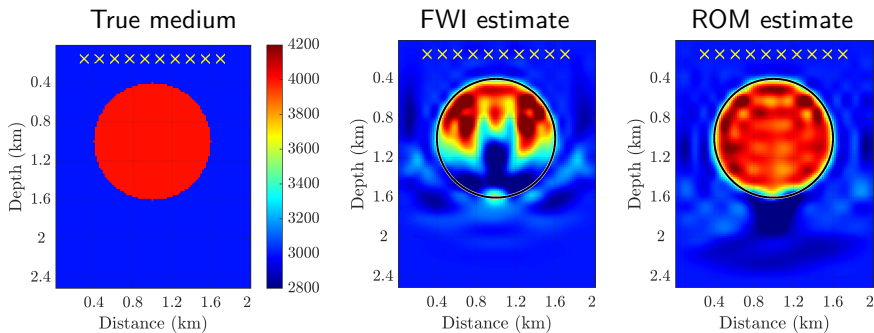
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- ROM objective function:

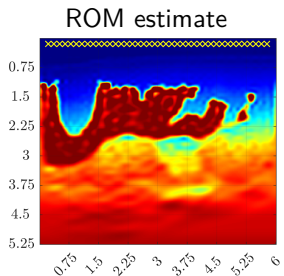
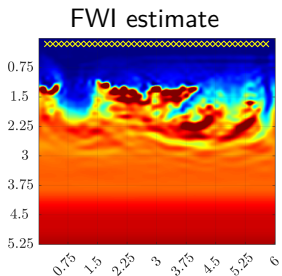
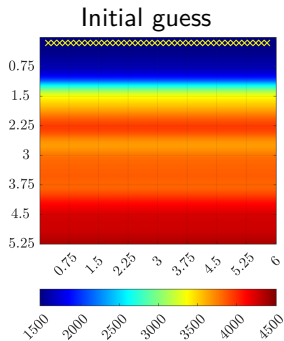
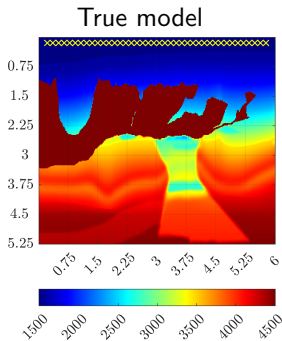
$$\mathcal{O}_{ROM}[c] = \|\mathbf{A}^{rom}[c] - \mathbf{A}^{rom,meas}\|_2^2$$

Camembert model



- Probing pulse is a modulated Gaussian pulse with central frequency 6Hz and bandwidth 4Hz ($\lambda = 300\text{m}$ at 10Hz).
- Search velocity: $c(x, \boldsymbol{\eta}) = c_o + \sum_I \eta_I \phi_I(x)$, $\boldsymbol{\eta} = (\eta_I)_{I=1}^L$.
- $\phi_I(x)$ are Gaussian peaks with centers on a regular grid, $L = 400$, with width 60m (0.2λ).
- FWI minimizes $\mathcal{O}_{FWI}(\boldsymbol{\eta}) = \|\mathcal{D}[c(\boldsymbol{\eta})] - \mathbf{d}^{meas}\|_2^2 + \mu \|\boldsymbol{\eta}\|_2^2$
- ROM minimizes $\mathcal{O}_{ROM}(\boldsymbol{\eta}) = \|\mathbf{A}^{rom}[c(\boldsymbol{\eta})] - \mathbf{A}^{rom,meas}\|_2^2 + \mu \|\boldsymbol{\eta}\|_2^2$

Salt body (BP - model)





A limitation and an extension to passive imaging


- One limitation of the method:
We need co-located sources and receivers.
- Extension to passive imaging:
Consider a receiver array recording signals transmitted by noise sources.
Compute the cross correlation matrix of the recorded signals.
 - The cross correlation matrix is related to the Green's function (virtual active array) [Shapiro et al. 2005; Garnier et al. 2016].
 - The ROM procedure is natural in the passive framework, since the cross correlation matrix gives directly the data matrix $\mathbf{D}(t)$.
 - the virtual sources and receivers are naturally co-located,
 - the signals are even (because cross correlations are even).

Conclusions

- The ROM is an approximation of the wave operator on a space defined by the snapshots of the wavefield at uniformly spaced time instants.
- This space is not known and neither is the wave operator.
- Yet, we can compute the ROM from the data !
- We can then formulate a velocity estimation algorithm that minimizes the ROM misfit and that avoids cycle skipping and other problems.
- Other formulations: See also

 L Borcea, J Garnier, AV Mamonov, J Zimmerling, When data driven reduced order modeling meets waveform inversion, arXiv:2302.05988

 L Borcea, J Garnier, AV Mamonov, J Zimmerling, Waveform inversion with a data driven estimate of the internal wave, SIAM Journal on Imaging Sciences 16 (1), 2023, 280-312.

 L Borcea, J Garnier, AV Mamonov, J Zimmerling, Waveform inversion via reduced order modeling, Geophysics 88 (2), 2023, R175-R191.